Newtonian Mechanics, Special Relativity and some General Relativity. Student Seminar Edvord Aksnes 2021-04-15

Dischaimen: I'm not a physicist, take everything with a lange grain of solt...

Questio: How can you describe what goes on in the miverse / our reality? Want to predict where "Things" are going. We can measure: · Space min ~ K • Time O Seems we live in $\mathbb{R} \times \mathbb{R}^3$.

Wamt: How do objects more in spacetime? Model: A big object - Single point. A motion: $\vec{x}: R \to R^3$ Worldhie.

Newtonian mechanics -> 2 physical concepts · Mass M · · F 1; F+F' 1 · Force Defined in vague Terms. Moss = "how much force to more something" Force = "smelling which makes sanething more"

Newton's laves: · 1st. law, principle of inentia: An object in motion stays in motion unless acted upon by a met external force $\sum F=0 \iff \frac{dV}{dF}=0$ · 2 nd law: Time change of momentum is proportional to the force $F = \frac{d \tilde{P}}{d t}$ p=mV monenton · 3 d Case: All forces between two objects exist in equal magnitude and opposite direction: FA = -FB

How are forces determined? Experiment Take an object, define its mass = 1. Measure have it mores under different contexts. Example: (Newton's law of universal gravitation) $|F_1| = |F_2| = G_1 \underbrace{m_1 m_2}_{q_{\text{partitude}} = G_1 \underbrace{m_1 m_2}_{\mathcal{R}^2}$ $\begin{pmatrix} m_1 \\ \vdots \end{pmatrix} \xrightarrow{F_1} \xrightarrow{F_2} (m_2)$ \hbar Modern theory: 4 Finder entel • Conavitation • Electromognetism • Strong force • Weak force More examples: FYS-MEK(100

How to change reference frame (coord. systen)? · Spacial notation: (X, t) (RX, t) REO(3) · Spacial translation: (\$\vec{x},t) > (\$\vec{x}+a,t) a \vec{R}^3 • Uniform motion: $(\vec{x},t) \mapsto (\vec{x}+t\vec{v},t) \vee eR^3$ Chalilean transformation: Corposition of these 3.

1860's: Maxwell's equation electromagnetism **Differential equations** Name Integral equations $\nabla \cdot \mathbf{E} = 4\pi \rho$ Gauss's law $\oint_{\partial \Omega} \mathbf{B} \cdot \mathrm{d}\mathbf{S} = 0$ Gauss's law for $\nabla \cdot \mathbf{B} = 0$ magnetism Maxwell-Faradav equation $abla imes {f E} = -rac{1}{c}rac{\partial {f B}}{\partial t}$ $\oint_{\partial \Sigma} \mathbf{E} \cdot \mathrm{d}\boldsymbol{\ell} = -\frac{1}{c} \frac{\mathrm{d}}{\mathrm{d}t} \iint_{\Sigma} \mathbf{B} \cdot \mathrm{d}\mathbf{S}$ (Faraday's law of induction) Ampère's circuital law $\oint_{\partial \Sigma} \mathbf{B} \cdot \mathrm{d}\boldsymbol{\ell} = \frac{1}{c} \left(4\pi \iint_{\Sigma} \mathbf{J} \cdot \mathrm{d}\mathbf{S} + \frac{\mathrm{d}}{\mathrm{d}t} \iint_{\Sigma} \mathbf{E} \cdot \mathrm{d}\mathbf{S} \right) \left| \nabla \times \mathbf{B} = \frac{1}{c} \left(4\pi \mathbf{J} + \frac{\partial \mathbf{E}}{\partial t} \right) \right|$ James Clerk Maxwell (with Maxwell's addition) $\mathbf{F} = q(\mathbf{E} + \mathbf{v} imes \mathbf{B})$ Losentz frice in variant under Gralikan transformation

However: Invariant under Lorentz transformation: $\begin{pmatrix} t \\ x \\ y \\ z \end{pmatrix} \mapsto \begin{pmatrix} \delta \left(t - \frac{y}{c^2} \right) \\ \delta \left(z - u \epsilon \right) \\ y \end{pmatrix}$ Velocity: $\vec{v} = (u, 0, 0)$ $V = \frac{1}{\sqrt{1 - \frac{v^2}{2}}}$ Lorent 2 factor. Observations: • 1. Time dilation 2. No simultaneity. $\{t=0\} \neq \{t'=0\} \subseteq \mathbb{R}^{+}$

Einstein : Galileon transformations Poiraré transformations (Lorentz tr. + translation Albert Einstein past Space Time IR time

Spocetime / Minkowski Space 4 manifold $\mathbb{R}^{1,3} = (\mathbb{R}^4, d)$ where $d((\epsilon_1, \vec{x}_1), (t_2, \vec{x}_2)) = \sqrt{-c^2(t_1-c_2)^2 + \|\vec{x}_1 - \vec{x}_2\|^2}$ $MeVnic: h_{\mu r} = \begin{pmatrix} -c^{2} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$ local coords: $(x^{\mu})_{\mu=0,\dots,3}$ timelike mull spoulike A porticle with moss follows a woldline where all tangent rectors are Einelike. Z: [01] -> R', 3, poor 2 lightcone $\Delta z = \int_{0}^{\infty} dz = \int_{0}^{\infty} \sqrt{\frac{1}{2}} - h_{\mu\nu} \frac{dz^{\nu}}{dz} \frac{dz^{\nu}}{dz} dz$

Given a particle with mass, we can parametrize its path with the internal dock AZ. $\|U\| = c^2$ $U = \left(\frac{d\widetilde{x}^{\mu}}{d\tau}\right)_{\mu=0,\dots,3}$ ~ Four - velocity no Momentum form-vector p = mU m mass of particle. Emergy is the 0-th component of p. $E = p^0$. In the reference france of the particle $U = (c^2, 0, 0, 0)$ $=> E = m c^2$

Suppose we have a reference france (= ") = 0, ..., 3, and a particle moving with relacity $v = \frac{dx}{dt}$ along to dr the x axis. In particle ref france: $(t', x') = (\tau, 0)$ $p = (mc^2, 0, 0, 0)$ In our ref frame: (t, x) Apply levent 2 transform $x = \delta(x' + vt') = \delta vt'$ P = (Y - Z) $\chi = \sqrt{1 - \frac{1}{2}}$ $\rho = \left(\chi_m c^2, \pi \chi_m, 0, 0 \right)$ Small v ~> p ~ m + 2 mv 2 (near energy + potential), p = mv (New lonion momentum)

Gravity is special Gravitational Mass. Inertial mass $\left[F_{g}=-m_{g}\nabla\phi\right]$ F = m; aWeak Equivolence Principle: $M_i = M_q$ Experimentally verified

ns Einstein Equivalence Principle In small enough negions of spocetime, the laves of physics reduce to those of special relativity. It is impossible to detect the existence of a gravitational field. Moving Rockets VS Tower t=0 t=1 $\begin{array}{ccc} t = 0 & a & a \\ \hline & \hline & \rightarrow & \hline & \rightarrow \end{array} \end{array}$ 173 ·m> t=1 $\square \rightarrow \square \rightarrow \\ \checkmark \rightarrow$

~ Spacetime is not R's it must be curred! In 1915, having worked for 8 years, Einstein published the Greneral theory of relativity. Curvature of spacetime and Matter & Energy Einstein's field equation: Rpv - 1 R gpv = 85tGrTpv Free particles: Greatisic equation: $\sum_{p,\sigma} \frac{d^2 x^p}{d \lambda^2} + \frac{\Gamma_p^p}{\sigma \frac{d x^p}{\delta \lambda}} \frac{d x^{\sigma}}{d \lambda^2} = 0$

Thanks for your time!

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